

$$e^t v'' + [2e^t + t(1-t)^{-1}e^t]v' = 2(1-t)e^{-t},$$

that is, $v'' + [(2-t)/(1-t)]v' = 2(1-t)e^{-2t}$. This equation is first order linear in v' , with integrating factor $\mu = e^t/(t-1)$. The solution is

$$v' = (t-1)(2e^{-2t} + c_1 e^{-t}).$$

Integrating, we obtain $v(t) = (1/2 - t)e^{-2t} - c_1 t e^{-t} + c_2$. Hence the solution of the original ODE is $y(t) = (1/2 - t)e^{-t} - c_1 t + c_2 e^t$.

Section 3.8

1. $R \cos \delta = 3$ and $R \sin \delta = 4 \Rightarrow R = \sqrt{25} = 5$ and $\delta = \arctan(4/3)$. Hence

$$u = 5 \cos(2t - 0.9273).$$

3. $R \cos \delta = 4$ and $R \sin \delta = -2 \Rightarrow R = \sqrt{20} = 2\sqrt{5}$ and $\delta = -\arctan(1/2)$. Hence

$$u = 2\sqrt{5} \cos(3t + 0.4636).$$

4. $R \cos \delta = -2$ and $R \sin \delta = -3 \Rightarrow R = \sqrt{13}$ and $\delta = \pi + \arctan(3/2)$. Hence

$$u = \sqrt{13} \cos(\pi t - 4.1244).$$

5. The spring constant is $k = 2/(1/2) = 4 \text{ lb/ft}$. Mass $m = 2/32 = 1/16 \text{ lb-s}^2/\text{ft}$. Since there is no damping, the equation of motion is

$$\frac{1}{16}u'' + 4u = 0,$$

that is, $u'' + 64u = 0$. The initial conditions are $u(0) = 1/4 \text{ ft}$, $u'(0) = 0 \text{ fps}$. The general solution is $u(t) = A \cos 8t + B \sin 8t$. Invoking the initial conditions, we have $u(t) = \frac{1}{4} \cos 8t$. $R = 3 \text{ inches}$, $\delta = 0 \text{ rad}$, $\omega_0 = 8 \text{ rad/s}$, and $T = \pi/4 \text{ sec}$.

7. The spring constant is $k = 3/(1/4) = 12 \text{ lb/ft}$. Mass $m = 3/32 \text{ lb-s}^2/\text{ft}$. Since there is no damping, the equation of motion is

$$\frac{3}{32}u'' + 12u = 0,$$

that is, $u'' + 128u = 0$. The initial conditions are $u(0) = -1/12 \text{ ft}$, $u'(0) = 2 \text{ fps}$. The general solution is $u(t) = A \cos 8\sqrt{2}t + B \sin 8\sqrt{2}t$. Invoking the initial conditions, we have

$$u(t) = -\frac{1}{12} \cos 8\sqrt{2}t + \frac{1}{4\sqrt{2}} \sin 8\sqrt{2}t.$$

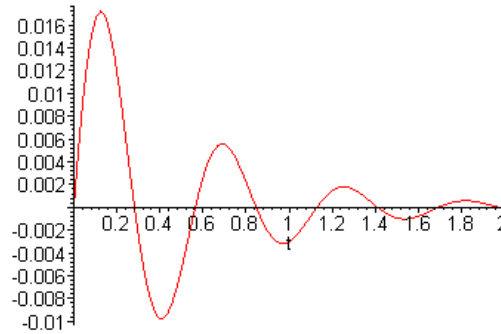
$R = \sqrt{11}/12 \text{ ft}$, $\delta = \pi - \text{atan}(3/\sqrt{2}) \text{ rad}$, $\omega_0 = 8\sqrt{2} \text{ rad/s}$, and $T = \pi/(4\sqrt{2}) \text{ sec}$.

10. The spring constant is $k = 16/(1/4) = 64 \text{ lb/ft}$. Mass $m = 1/2 \text{ lb-s}^2/\text{ft}$. The damping coefficient is $\gamma = 2 \text{ lb-sec/ft}$. Hence the equation of motion is

$$\frac{1}{2}u'' + 2u' + 64u = 0,$$

that is, $u'' + 4u' + 128u = 0$. The initial conditions are $u(0) = 0 \text{ ft}$, $u'(0) = 1/4 \text{ fps}$. The general solution is $u(t) = A \cos 2\sqrt{31}t + B \sin 2\sqrt{31}t$. Invoking the initial conditions, we have

$$u(t) = \frac{1}{8\sqrt{31}} e^{-2t} \sin 2\sqrt{31}t.$$



Solving $u(t) = 0$, on the interval $[0.2, 0.4]$, we obtain $t = \pi/2\sqrt{31} = 0.2821 \text{ sec}$. Based on the graph, and the solution of $u(t) = 0.01$, we have $|u(t)| \leq 0.01$ for $t \geq \tau = 0.2145$.

11. The spring constant is $k = 3/(.1) = 30 \text{ N/m}$. The damping coefficient is given as $\gamma = 3/5 \text{ N-sec/m}$. Hence the equation of motion is

$$2u'' + \frac{3}{5}u' + 30u = 0,$$

that is, $u'' + 0.3u' + 15u = 0$. The initial conditions are $u(0) = 0.05 \text{ m}$ and $u'(0) = 0.01 \text{ m/s}$. The general solution is $u(t) = A \cos \mu t + B \sin \mu t$, in which $\mu = 3.87008 \text{ rad/s}$. Invoking the initial conditions, we have

$$u(t) = e^{-0.15t}(0.05 \cos \mu t + 0.00452 \sin \mu t).$$

Also, $\mu/\omega_0 = 3.87008/\sqrt{15} \approx 0.99925$.

13. The frequency of the *undamped* motion is $\omega_0 = 1$. The quasi frequency of the damped motion is $\mu = \frac{1}{2}\sqrt{4 - \gamma^2}$. Setting $\mu = \frac{2}{3}\omega_0$, we obtain $\gamma = \frac{2}{3}\sqrt{5}$.

14. The spring constant is $k = mg/L$. The equation of motion for an undamped system is

$$mu'' + \frac{mg}{L}u = 0.$$

Hence the natural frequency of the system is $\omega_0 = \sqrt{\frac{g}{L}}$. The period is $T = 2\pi/\omega_0$.

15. The general solution of the system is $u(t) = A \cos \gamma(t - t_0) + B \sin \gamma(t - t_0)$. Invoking the initial conditions, we have $u(t) = u_0 \cos \gamma(t - t_0) + (u'_0/\gamma) \sin \gamma(t - t_0)$. Clearly, the functions $v = u_0 \cos \gamma(t - t_0)$ and $w = (u'_0/\gamma) \sin \gamma(t - t_0)$ satisfy the given criteria.

16. Note that $r \sin(\omega_0 t - \theta) = r \sin \omega_0 t \cos \theta - r \cos \omega_0 t \sin \theta$. Comparing the given expressions, we have $A = -r \sin \theta$ and $B = r \cos \theta$. That is, $r = R = \sqrt{A^2 + B^2}$, and $\tan \theta = -A/B = -1/\tan \delta$. The latter relation is also $\tan \theta + \cot \delta = 1$.

18. The system is *critically damped*, when $R = 2\sqrt{L/C}$. Here $R = 1000$ ohms.

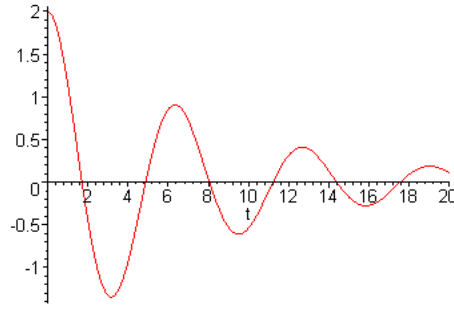
21(a). Let $u = Re^{-\gamma t/2m} \cos(\mu t - \delta)$. Then attains a *maximum* when $\mu t_k - \delta = 2k\pi$. Hence $T_d = t_{k+1} - t_k = 2\pi/\mu$.

(b). $u(t_k)/u(t_{k+1}) = \exp(-\gamma t_k/2m)/\exp(-\gamma t_{k+1}/2m) = \exp[(\gamma t_{k+1} - \gamma t_k)/2m]$. Hence $u(t_k)/u(t_{k+1}) = \exp[\gamma(2\pi/\mu)/2m] = \exp(\gamma T_d/2m)$.

(c). $\Delta = \ln[u(t_k)/u(t_{k+1})] = \gamma(2\pi/\mu)/2m = \pi\gamma/\mu m$.

22. The spring constant is $k = 16/(1/4) = 64$ lb/ft. Mass $m = 1/2$ lb-s²/ft. The damping coefficient is $\gamma = 2$ lb-sec/ft. The quasi frequency is $\mu = 2\sqrt{31}$ rad/s. Hence $\Delta = \frac{2\pi}{\sqrt{31}} \approx 1.1285$.

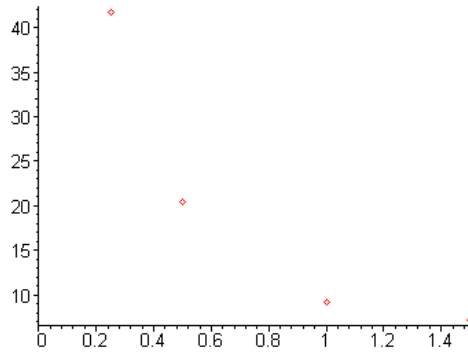
25(a). The solution of the IVP is $u(t) = e^{-t/8} \left(2 \cos \frac{3}{8} \sqrt{7} t + 0.252 \sin \frac{3}{8} \sqrt{7} t \right)$.



Using the plot, and numerical analysis, $\tau \approx 41.715$.

(b). For $\gamma = 0.5$, $\tau \approx 20.402$; for $\gamma = 1.0$, $\tau \approx 9.168$; for $\gamma = 1.5$, $\tau \approx 7.184$.

(c).



(d). For $\gamma = 1.6$, $\tau \approx 7.218$; for $\gamma = 1.7$, $\tau \approx 6.767$; for $\gamma = 1.8$, $\tau \approx 5.473$; for $\gamma = 1.9$, $\tau \approx 6.460$. τ steadily decreases to about $\tau_{min} \approx 4.873$, corresponding to the critical value $\gamma_0 \approx 1.73$.

(e). We have $u(t) = \frac{4e^{-\gamma t/2}}{\sqrt{4-\gamma^2}} \cos(\mu t - \delta)$, in which $\mu = \frac{1}{2}\sqrt{4-\gamma^2}$, and $\delta = \tan^{-1} \frac{\gamma}{\sqrt{4-\gamma^2}}$. Hence $|u(t)| \leq \frac{4e^{-\gamma t/2}}{\sqrt{4-\gamma^2}}$.

26(a). The characteristic equation is $mr^2 + \gamma r + k = 0$. Since $\gamma^2 < 4km$, the roots are $r_{1,2} = -\frac{\gamma}{2m} \pm i \frac{\sqrt{4mk-\gamma^2}}{2m}$. The general solution is

$$u(t) = e^{-\gamma t/2m} \left[A \cos \frac{\sqrt{4mk-\gamma^2}}{2m} t + B \sin \frac{\sqrt{4mk-\gamma^2}}{2m} t \right].$$

Invoking the initial conditions, $A = u_0$ and

$$B = \frac{(2mv_0 - \gamma u_0)}{\sqrt{4mk - \gamma^2}}.$$

(b). We can write $u(t) = R e^{-\gamma t/2m} \cos(\mu t - \delta)$, in which

$$R = \sqrt{u_0^2 + \frac{(2mv_0 - \gamma u_0)^2}{4mk - \gamma^2}},$$

and

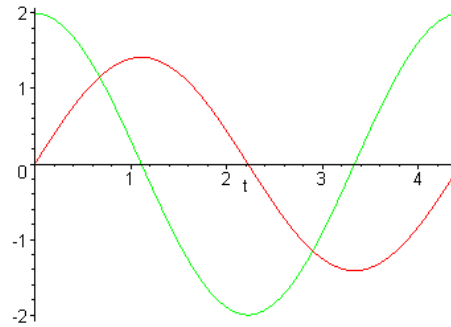
$$\delta = \arctan \left[\frac{(2mv_0 - \gamma u_0)}{u_0 \sqrt{4mk - \gamma^2}} \right].$$

(c). $R = \sqrt{u_0^2 + \frac{(2mv_0 - \gamma u_0)^2}{4mk - \gamma^2}} = 2\sqrt{\frac{m(ku_0^2 + \gamma u_0 v_0 + mv_0^2)}{4mk - \gamma^2}} = \sqrt{\frac{a+b\gamma}{4mk - \gamma^2}}.$

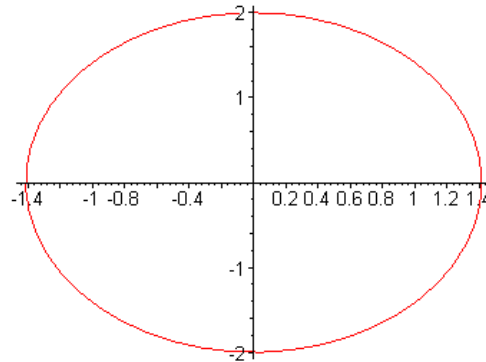
It is evident that R increases (*monotonically*) without bound as $\gamma \rightarrow (2\sqrt{mk})^-$.

28(a). The general solution is $u(t) = A \cos \sqrt{2}t + B \sin \sqrt{2}t$. Invoking the initial conditions, we have $u(t) = \sqrt{2} \sin \sqrt{2}t$.

(b).

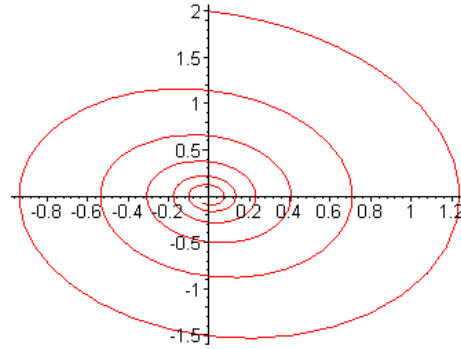


(c).



The condition $u'(0) = 2$ implies that $u(t)$ *initially* increases. Hence the phase point travels *clockwise*.

29. $u(t) = \frac{16}{\sqrt{127}} e^{-t/8} \sin \frac{\sqrt{127}}{8} t.$



31. Based on *Newton's second law*, with the positive direction to the right,

$$\sum F = mu''$$

where

$$\sum F = -ku - \gamma u'.$$

Hence the equation of motion is $mu'' + \gamma u' + ku = 0$. The only difference in this problem is that the equilibrium position is located at the *unstretched* configuration of the spring.

32(a). The *restoring* force exerted by the spring is $F_s = -(ku + \varepsilon u^3)$. The *opposing* viscous force is $F_d = -\gamma u'$. Based on *Newton's second law*, with the positive direction to the right,

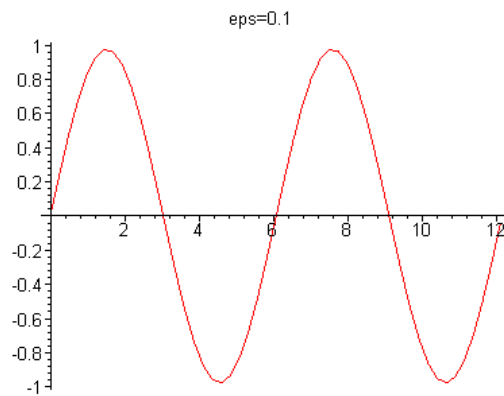
$$F_s + F_d = mu''.$$

Hence the equation of motion is $mu'' + \gamma u' + ku + \varepsilon u^3 = 0$.

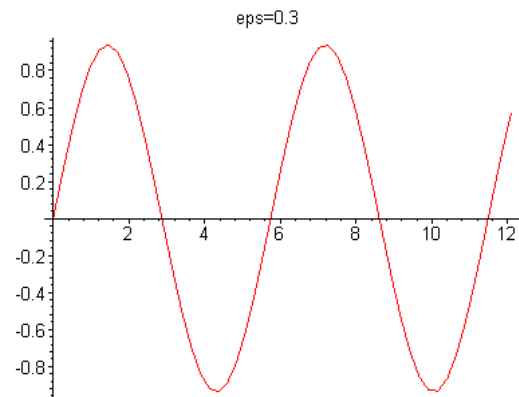
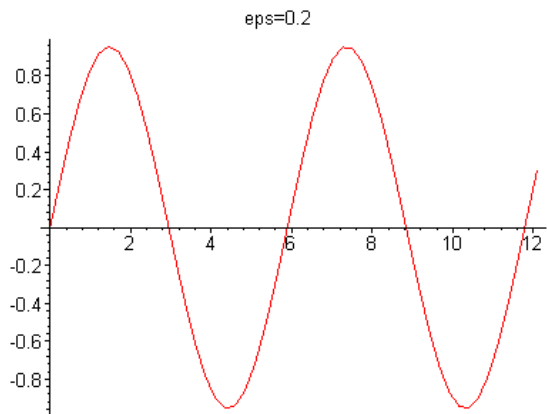
(b). With the specified parameter values, the equation of motion is $u'' + u = 0$. The general solution of this ODE is $u(t) = A \cos t + B \sin t$. Invoking the initial conditions, the specific solution is $u(t) = \sin t$. Clearly, the amplitude is $R = 1$, and the period of the motion is $T = 2\pi$.

(c). Given $\varepsilon = 0.1$, the equation of motion is $u'' + u + 0.1 u^3 = 0$. A solution of the

IVP can be generated numerically:



(d).



(e). The amplitude and period both seem to *decrease*.

(f).

