

## Section 9.2

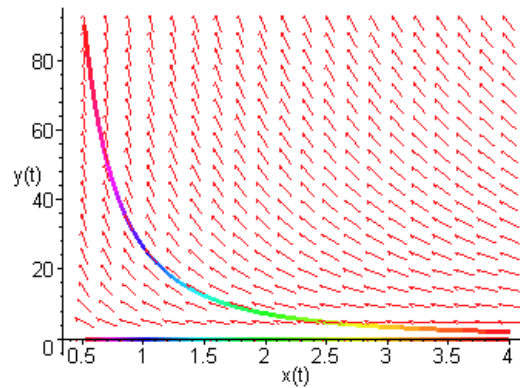
2. The differential equations can be combined to obtain a related ODE

$$\frac{dy}{dx} = -\frac{2y}{x}.$$

The equation is *separable*, with

$$\frac{dy}{y} = -\frac{2 dx}{x}.$$

The solution is given by  $y = C x^{-2}$ . Note that the system is *uncoupled*, and hence we also have  $x = x_0 e^{-t}$  and  $y = y_0 e^{2t}$ .



In order to determine the direction of motion along the trajectories, observe that for *positive* initial conditions,  $x$  will *decrease*, whereas  $y$  will *increase*.

4. The trajectories of the system satisfy the ODE

$$\frac{dy}{dx} = -\frac{bx}{ay}.$$

The equation is *separable*, with

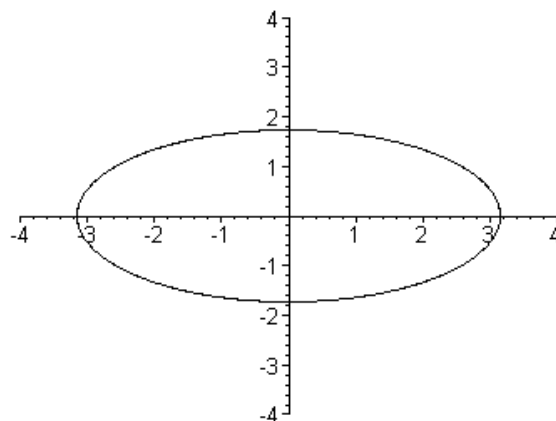
$$ay dy = -bx dx.$$

Hence the trajectories are given by  $b x^2 + a y^2 = C^2$ , in which  $C$  is arbitrary. Evidently, the trajectories are *ellipses*. Invoking the initial condition, we find that  $C^2 = ab$ . The system of ODEs can also be written as

$$\frac{d\mathbf{x}}{dt} = \begin{pmatrix} 0 & a \\ -b & 0 \end{pmatrix} \mathbf{x}.$$

Using the methods in Chapter 7, it is easy to show that

$$\begin{aligned} x &= \sqrt{a} \cos \sqrt{ab} t \\ y &= -\sqrt{b} \sin \sqrt{ab} t. \end{aligned}$$



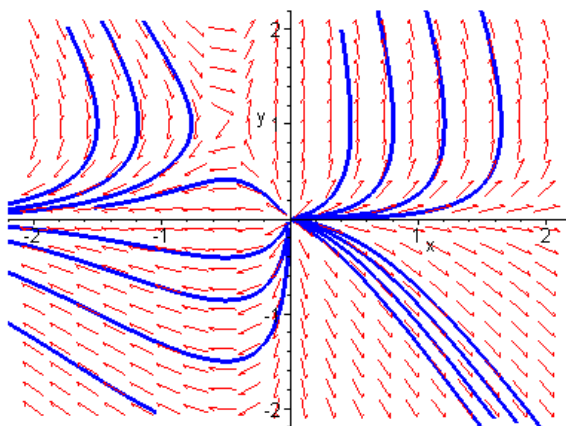
Note that for *positive* initial conditions,  $x$  will *increase*, whereas  $y$  will *decrease*.

5(a). The critical points are given by the solution set of the equations

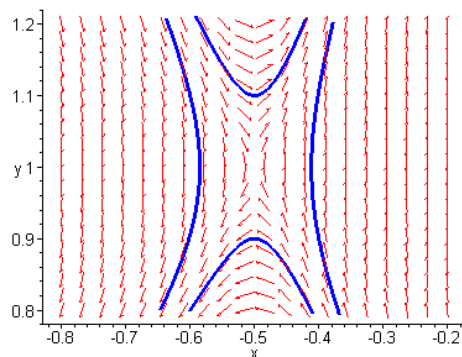
$$\begin{aligned} x(1 - y) &= 0 \\ y(1 + 2x) &= 0. \end{aligned}$$

Clearly,  $(0, 0)$  is a solution. If  $x \neq 0$ , then  $y = 1$  and  $x = -1/2$ . Hence the critical points are  $(0, 0)$  and  $(-1/2, 1)$ .

(b).



(c). Based on the phase portrait, all trajectories starting near the origin *diverge*. Hence the critical point  $(0, 0)$  is *unstable*. Examining the phase curves near the critical point  $(-1/2, 1)$ ,



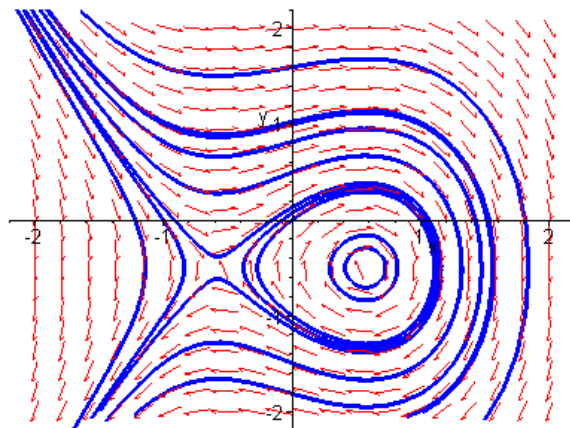
the equilibrium point has the properties of a *saddle*, and hence it is *unstable*.

6(a). The critical points are solutions of the equations

$$\begin{aligned} 1 + 2y &= 0 \\ 1 - 3x^2 &= 0. \end{aligned}$$

There are two equilibrium points,  $\left(-1/\sqrt{3}, -1/2\right)$  and  $\left(1/\sqrt{3}, -1/2\right)$ .

(b).



(c). Locally, the trajectories near the point  $\left(-1/\sqrt{3}, -1/2\right)$  resemble the behavior near a *saddle*. Hence the critical point is *unstable*. Near the point  $\left(1/\sqrt{3}, -1/2\right)$ , the solutions are *periodic*. Therefore the second critical point is *stable*.

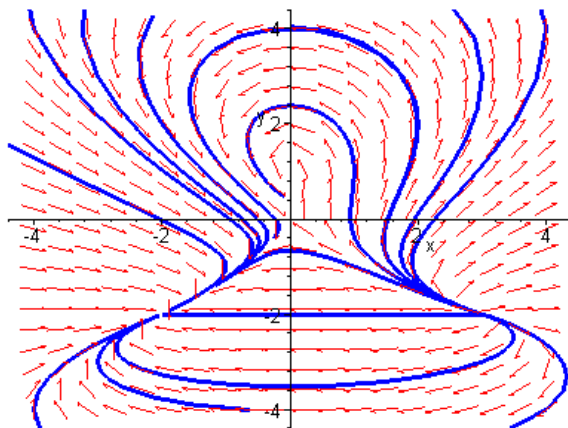
8(a). The critical points are solutions of the equations

$$\begin{aligned} -(x - y)(1 - x - y) &= 0 \\ x(2 + y) &= 0. \end{aligned}$$

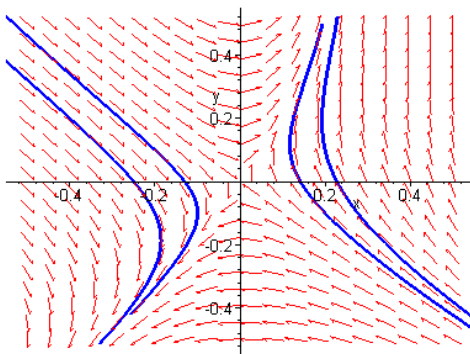
If  $x = y$ , then  $x = y = 0$  or  $x = y = -2$ . If  $x = 1 - y$ , then  $x = 0$  and  $y = 1$ , or  $x = 3$  and  $y = -2$ . It follows that the critical points are  $(0, 0)$ ,  $(-2, -2)$ ,  $(0, 1)$

and  $(3, -2)$ .

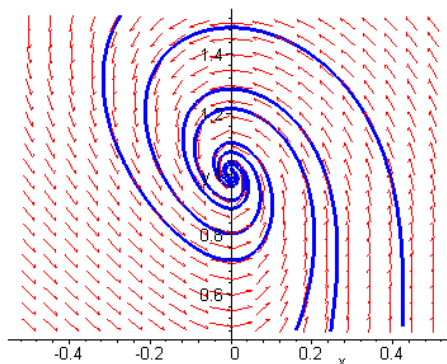
(b).



(c). Near the origin, the trajectories resemble those of a *saddle*, and hence it is *unstable*.



Near the critical point  $(0, 1)$ , the trajectories resemble those of a stable *spiral*. Hence the equilibrium point is *asymptotically stable*.



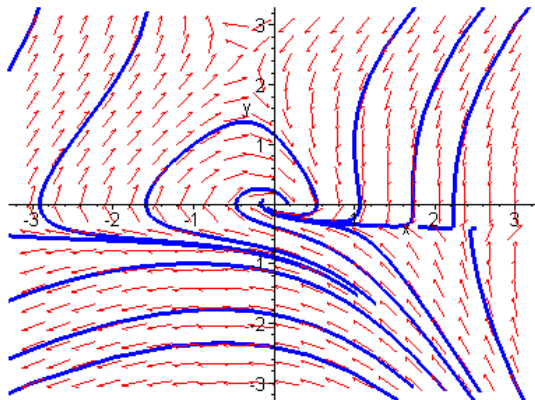
Based on the global phase portrait, it is evident that the other critical points are *nodes*. Closer examination reveals that the point  $(-2, -2)$  is *asymptotically stable*, whereas the point  $(3, -2)$  is *unstable*.

9(a). The critical points are given by the solution set of the equations

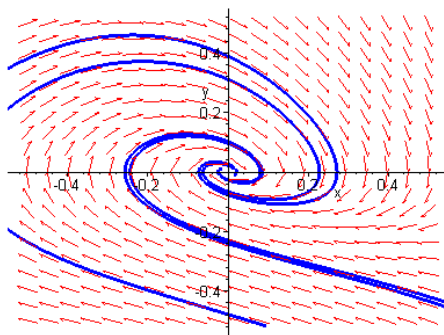
$$\begin{aligned} y(2 - x - y) &= 0 \\ -x - y - 2xy &= 0. \end{aligned}$$

Clearly,  $(0, 0)$  is a critical point. If  $x = 2 - y$ , then it follows that  $y(y - 2) = 1$ . The additional critical points are  $(1 - \sqrt{2}, 1 + \sqrt{2})$  and  $(1 + \sqrt{2}, 1 - \sqrt{2})$ .

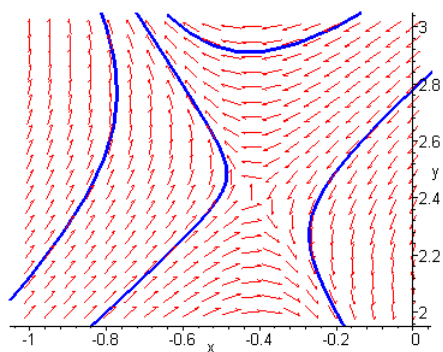
(b).



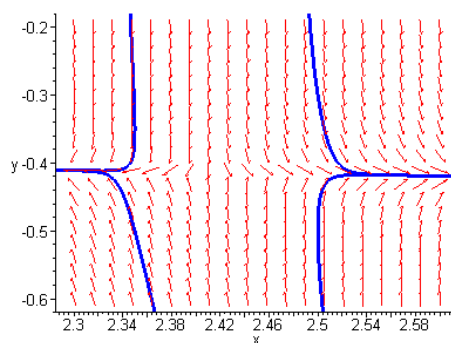
(c). The behavior near the origin is that of a *stable spiral*. Hence the point  $(0, 0)$  is *asymptotically stable*.



At the critical point  $(1 - \sqrt{2}, 1 + \sqrt{2})$ , the trajectories resemble those near a *saddle*. Hence the critical point is *unstable*.



Near the point  $(1 + \sqrt{2}, 1 - \sqrt{2})$ , the trajectories resemble those near a *saddle*.  
Hence the critical point is also *unstable*.

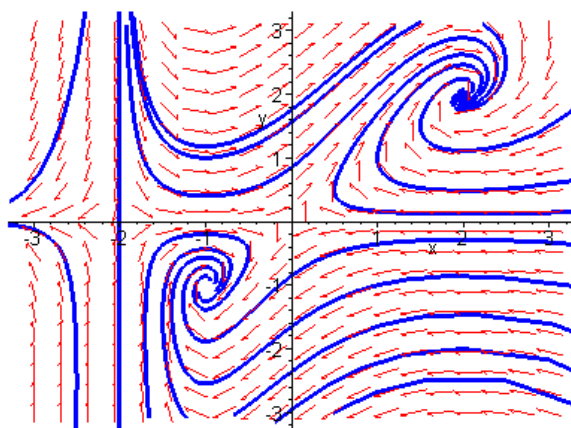


10(a). The critical points are solutions of the equations

$$\begin{aligned}(2+x)(y-x) &= 0 \\ y(2+x-x^2) &= 0.\end{aligned}$$

The origin is evidently a critical point. If  $x = -2$ , then  $y = 0$ . If  $x = y$ , then either  $y = 0$  or  $x = y = -1$  or  $x = y = 2$ . Hence the other critical points are  $(-2, 0)$ ,  $(-1, -1)$  and  $(2, 2)$ .

(b).



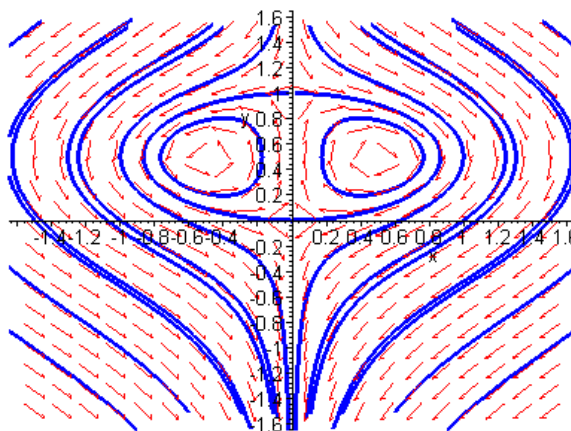
(c). Based on the global phase portrait, the critical points  $(0, 0)$  and  $(-2, 0)$  have the characteristics of a *saddle*. Hence these points are *unstable*. The behavior near the remaining two critical points resembles those near a *stable spiral*. Hence the critical points  $(-1, -1)$  and  $(2, 2)$  are *asymptotically stable*.

11(a). The critical points are given by the solution set of the equations

$$\begin{aligned} x(1 - 2y) &= 0 \\ y - x^2 - y^2 &= 0. \end{aligned}$$

If  $x = 0$ , then either  $y = 0$  or  $y = 1$ . If  $y = 1/2$ , then  $x = \pm 1/2$ . Hence the critical points are at  $(0, 0)$ ,  $(0, 1)$ ,  $(-1/2, 1/2)$  and  $(1/2, 1/2)$ .

(b).



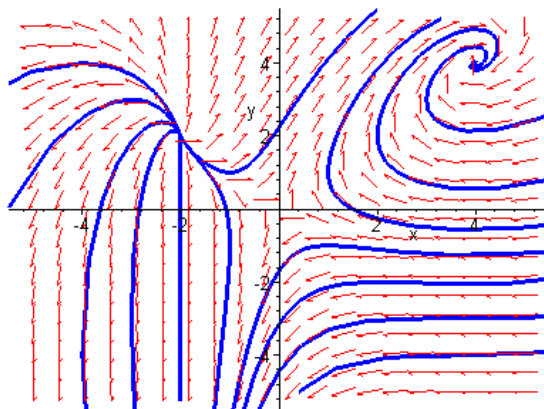
(c). The trajectories near the critical points  $(-1/2, 1/2)$  and  $(1/2, 1/2)$  are closed curves. Hence the critical points have the characteristics of a *center*, which is *stable*. The trajectories near the critical points  $(0, 0)$  and  $(0, 1)$  resemble those near a *saddle*. Hence these critical points are *unstable*.

13(a). The critical points are solutions of the equations

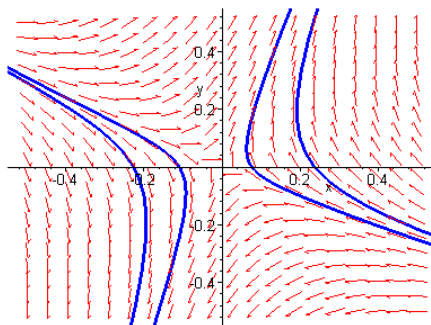
$$\begin{aligned}(2+x)(y-x) &= 0 \\ (4-x)(y+x) &= 0.\end{aligned}$$

If  $y = x$ , then either  $x = y = 0$  or  $x = y = 4$ . If  $x = -2$ , then  $y = 2$ . If  $x = -y$ , then  $y = 2$  or  $y = 0$ . Hence the critical points are at  $(0, 0)$ ,  $(4, 4)$  and  $(-2, 2)$ .

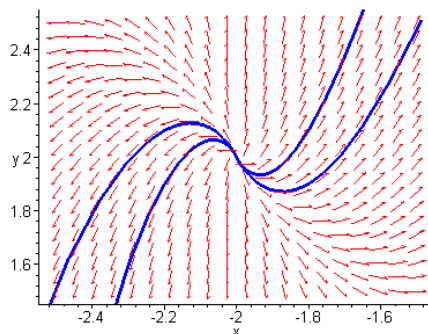
(b).



(c). The critical point at  $(4, 4)$  is evidently a *stable spiral*, which is *asymptotically stable*. Closer examination of the critical point at  $(0, 0)$  reveals that it is a *saddle*, which is *unstable*.



The trajectories near the critical point  $(-2, 2)$  resemble those near an *unstable node*.

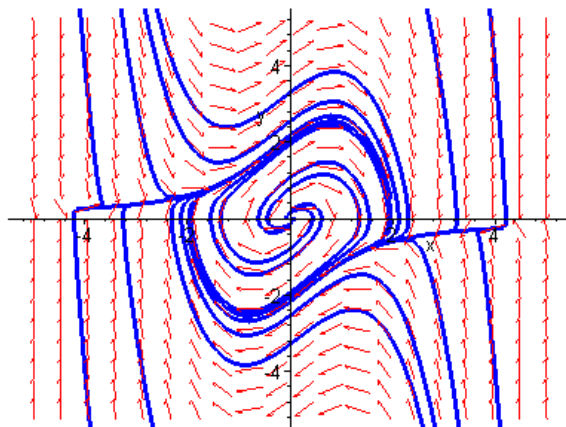


14(a). The critical points consist of the solution set of the equations

$$\begin{aligned} y &= 0 \\ (1 - x^2)y - x &= 0. \end{aligned}$$

It is easy to see that the only critical point is at  $(0, 0)$ .

(b).



(c). The origin is an *unstable spiral*.

16(a). The trajectories are solutions of the differential equation

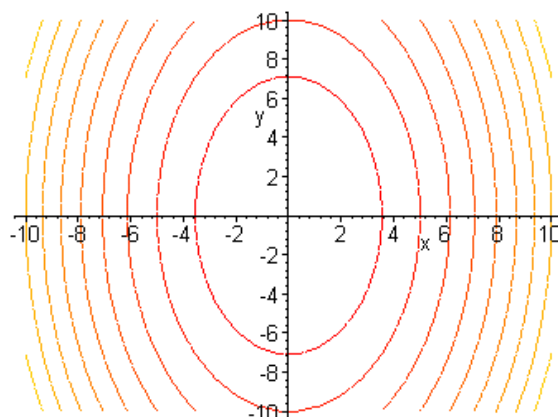
$$\frac{dy}{dx} = -\frac{4x}{y},$$

which can also be written as  $4x \, dx + y \, dy = 0$ . Integrating, we obtain

$$4x^2 + y^2 = C^2.$$

Hence the trajectories are ellipses.

(b).



Based on the differential equations, the direction of motion on each trajectory is *clockwise*.

17(a). The trajectories of the system satisfy the ODE

$$\frac{dy}{dx} = \frac{2x + y}{y},$$

which can also be written as  $(2x + y)dx - ydy = 0$ . This differential equation is *homogeneous*. Setting  $y = xv(x)$ , we obtain

$$v + x \frac{dv}{dx} = \frac{2}{v} + 1,$$

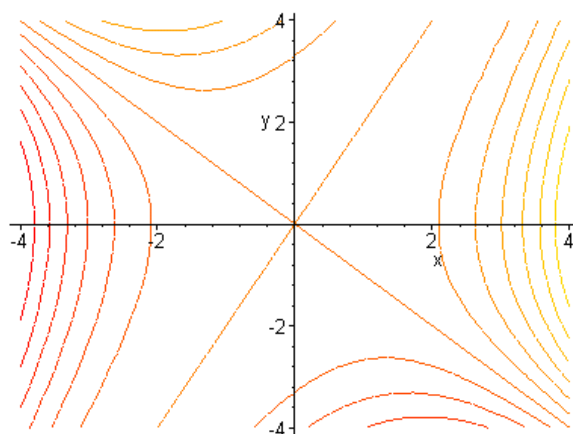
that is,

$$x \frac{dv}{dx} = \frac{2 + v - v^2}{v}.$$

The resulting ODE is *separable*, with solution  $x^3(v + 1)(v - 2)^2 = C$ . Reverting back to the original variables, the trajectories are level curves of

$$H(x, y) = (x + y)(y - 2x)^2.$$

(b).



The origin is a *saddle*. Along the line  $y = 2x$ , solutions increase without bound. Along the line  $y = -x$ , solutions converge toward the origin.

18(a). The trajectories are solutions of the differential equation

$$\frac{dy}{dx} = \frac{x+y}{x-y},$$

which is *homogeneous*. Setting  $y = x v(x)$ , we obtain

$$v + x \frac{dv}{dx} = \frac{x + xv}{x - xv},$$

that is,

$$x \frac{dv}{dx} = \frac{1 + v^2}{1 - v}.$$

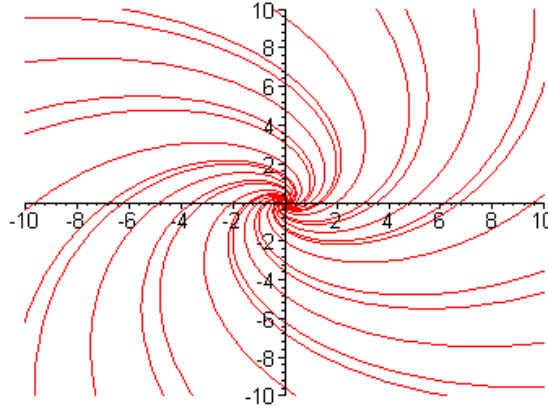
The resulting ODE is *separable*, with solution

$$\arctan(v) = \ln|x| \sqrt{1 + v^2}.$$

Reverting back to the original variables, the trajectories are level curves of

$$H(x, y) = \arctan(y/x) - \ln \sqrt{x^2 + y^2}.$$

(b).



The origin is a *stable spiral*.

20(a). The trajectories are solutions of the differential equation

$$\frac{dy}{dx} = \frac{-2xy^2 + 6xy}{2x^2y - 3x^2 - 4y},$$

which can also be written as  $(2xy^2 - 6xy)dx + (2x^2y - 3x^2 - 4y)dy = 0$ . The resulting ODE is *exact*, with

$$\frac{\partial H}{\partial x} = 2xy^2 - 6xy \quad \text{and} \quad \frac{\partial H}{\partial y} = 2x^2y - 3x^2 - 4y.$$

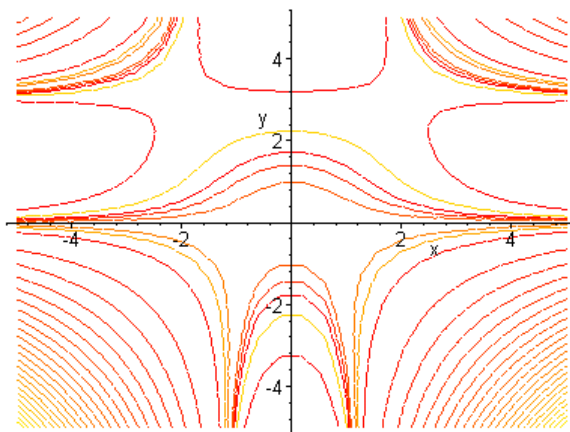
Integrating the first equation, we find that  $H(x, y) = x^2y^2 - 3x^2y + f(y)$ . It follows that

$$\frac{\partial H}{\partial y} = 2x^2y - 3x^2 + f'(y).$$

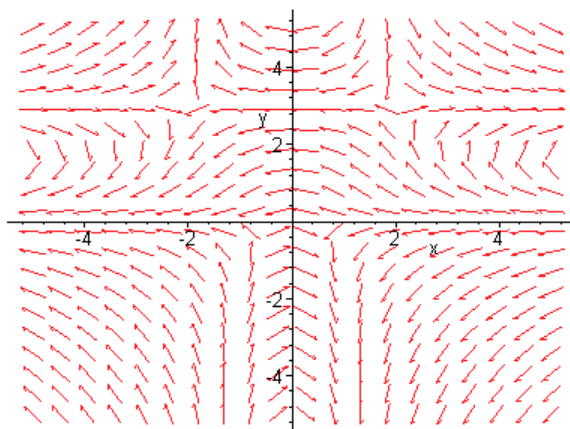
Comparing the two partial derivatives, we obtain  $f(y) = -2y^2 + c$ . Hence

$$H(x, y) = x^2y^2 - 3x^2y - 2y^2.$$

(b).



The associated direction field shows the direction of motion along the trajectories.



22(a). The trajectories are solutions of the differential equation

$$\frac{dy}{dx} = \frac{-6x + x^3}{6y},$$

which can also be written as  $(6x - x^3)dx + 6ydy = 0$ . The resulting ODE is *exact*, with

$$\frac{\partial H}{\partial x} = 6x - x^3 \text{ and } \frac{\partial H}{\partial y} = 6y.$$

Integrating the first equation, we have  $H(x, y) = 3x^2 - x^4/4 + f(y)$ . It follows that

$$\frac{\partial H}{\partial y} = f'(y).$$

Comparing the two partial derivatives, we conclude that  $f(y) = 3y^2 + c$ . Hence

$$H(x, y) = 3x^2 - \frac{x^4}{4} + 3y^2.$$

(b).

