

Section 11.4

1. Let $\phi_n(x) = J_0(\sqrt{\lambda_n} x)$ be the eigenfunctions of the singular problem

$$\begin{aligned} -(xy')' &= \lambda xy, \quad 0 < x < 1, \\ y, y' &\text{ bounded as } x \rightarrow 0, \quad y(1) = 0. \end{aligned}$$

Let $\phi(x)$ be a solution of the given BVP, and set

$$\phi(x) = \sum_{n=0}^{\infty} b_n \phi_n(x). \quad (*)$$

Then

$$\begin{aligned} -(x\phi')' &= \mu x\phi + f(x) \\ &= \mu x\phi + x \frac{f(x)}{x}. \end{aligned}$$

Substituting $(*)$, we obtain

$$\sum_{n=0}^{\infty} b_n \lambda_n x \phi_n(x) = \mu x \sum_{n=0}^{\infty} b_n \phi_n(x) + x \sum_{n=0}^{\infty} c_n \phi_n(x),$$

in which the c_n are the expansion coefficients of $f(x)/x$ for $x > 0$. That is,

$$\begin{aligned} c_n &= \frac{1}{\|\phi_n(x)\|^2} \int_0^1 x \frac{f(x)}{x} \phi_n(x) dx \\ &= \frac{1}{\|\phi_n(x)\|^2} \int_0^1 f(x) \phi_n(x) dx. \end{aligned}$$

It follows that if $x \neq 0$,

$$\sum_{n=0}^{\infty} [c_n - b_n(\lambda_n - \mu)] \phi_n(x) = 0.$$

As long as $\mu \neq \lambda_n$, linear independence of the eigenfunctions implies that

$$b_n = \frac{c_n}{\lambda_n - \mu}, \quad n = 1, 2, \dots.$$

Therefore a formal solution is given by

$$\phi(x) = \sum_{n=0}^{\infty} \frac{c_n}{\lambda_n - \mu} J_0(\sqrt{\lambda_n} x),$$

in which $\sqrt{\lambda_n}$ are the positive roots of $J_0(x) = 0$.

3(a). Setting $t = \sqrt{\lambda} x$, it follows that

$$\frac{dy}{dx} = \sqrt{\lambda} \frac{dy}{dt} \quad \text{and} \quad \frac{d^2y}{dx^2} = \lambda \frac{d^2y}{dt^2}.$$

The given ODE can be expressed as

$$-\sqrt{\lambda} \frac{d}{dt} \left(\frac{t}{\sqrt{\lambda}} \sqrt{\lambda} \frac{dy}{dt} \right) + \frac{k^2 \sqrt{\lambda}}{t} = \sqrt{\lambda} t y,$$

or

$$-\frac{d}{dt} \left(t \frac{dy}{dt} \right) + \frac{k^2}{t} = t y.$$

An equivalent form is given by

$$t^2 \frac{dy}{dt} + t \frac{dy}{dt} + (t^2 - k^2)y = 0,$$

which is known as a Bessel equation of order k . A *bounded* solution is $J_k(t)$.

(b). $J_k(\sqrt{\lambda} x)$ satisfies the boundary condition at $x = 0$. Imposing the other boundary

condition, it is necessary that $J_k(\sqrt{\lambda}) = 0$. Therefore the eigenvalues are given by λ_n , $n = 1, 2, \dots$, where $\sqrt{\lambda_n}$ are the positive zeroes of $J_k(x)$. The eigenfunctions of the BVP are $\phi_n(x) = J_k(\sqrt{\lambda_n} x)$.

(c). The BVP is a *singular Sturm-Liouville* problem with

$$L[y] = -(xy')' + \frac{k^2}{x} y \quad \text{and} \quad r(x) = 1.$$

We note that

$$\begin{aligned} \lambda_n \int_0^1 x \phi_n(x) \phi_m(x) dx &= \int_0^1 L[\phi_n] \phi_m(x) dx \\ &= \int_0^1 \phi_n(x) L[\phi_m] dx \\ &= \lambda_m \int_0^1 x \phi_n(x) \phi_m(x) dx. \end{aligned}$$

Therefore

$$(\lambda_n - \lambda_m) \int_0^1 x \phi_n(x) \phi_m(x) dx = 0.$$

So for $n \neq m$, we have $\lambda_n \neq \lambda_m$ and

$$\int_0^1 x \phi_n(x) \phi_m(x) dx = 0.$$

(d). Consider the expansion

$$f(x) = \sum_{n=0}^{\infty} a_n \phi_n(x).$$

Multiplying both sides of equation by $x \phi_j(x)$ and integrating from 0 to 1, and using the orthogonality of the eigenfunction,

$$\begin{aligned} \int_0^1 x f(x) \phi_j(x) dx &= \sum_{n=0}^{\infty} a_n \int_0^1 x \phi_j(x) \phi_n(x) dx \\ &= a_j \int_0^1 x \phi_j(x) \phi_j(x) dx. \end{aligned}$$

Therefore

$$a_j = \int_0^1 x f(x) \phi_j(x) dx / \int_0^1 x [\phi_j(x)]^2 dx, \quad j = 1, 2, \dots.$$

(e). Let $\phi(x)$ be a solution of the given BVP, and set

$$\phi(x) = \sum_{n=0}^{\infty} b_n \phi_n(x), \tag{*}$$

where $\phi_n(x) = J_k(\sqrt{\lambda_n} x)$. Then

$$\begin{aligned} L[\phi] &= \mu x \phi + f(x) \\ &= \mu x \phi + x \frac{f(x)}{x}. \end{aligned}$$

Substituting (*), we obtain

$$\sum_{n=0}^{\infty} b_n \lambda_n x \phi_n(x) = \mu x \sum_{n=0}^{\infty} b_n \phi_n(x) + x \sum_{n=0}^{\infty} c_n \phi_n(x),$$

in which the c_n are the expansion coefficients of $f(x)/x$ for $x > 0$. That is,

$$\begin{aligned} c_n &= \frac{1}{\|\phi_n(x)\|^2} \int_0^1 x \frac{f(x)}{x} \phi_n(x) dx \\ &= \frac{1}{\|J_k(\sqrt{\lambda_n} x)\|^2} \int_0^1 f(x) J_k(\sqrt{\lambda_n} x) dx. \end{aligned}$$

It follows that if $x \neq 0$,

$$\sum_{n=0}^{\infty} [c_n - b_n(\lambda_n - \mu)] J_k(\sqrt{\lambda_n} x) = 0.$$

As long as $\mu \neq \lambda_n$, linear independence of the eigenfunctions implies that

$$b_n = \frac{c_n}{\lambda_n - \mu}, \quad n = 1, 2, \dots.$$

Therefore a formal solution is given by

$$\phi(x) = \sum_{n=0}^{\infty} \frac{c_n}{\lambda_n - \mu} J_k(\sqrt{\lambda_n} x).$$

5(a). Setting $\lambda = \alpha^2$ in Prob. 15 of Section 11.1, the *Chebyshev equation* can also be written as

$$-\left[\sqrt{1-x^2} y'\right]' = \frac{\lambda}{\sqrt{1-x^2}} y.$$

Note that

$$p(x) = \sqrt{1-x^2}, \quad q(x) = 0, \quad \text{and} \quad r(x) = 1/\sqrt{1-x^2},$$

hence both boundary points are singular.

(b). Observe that $p(1-\varepsilon) = \sqrt{2\varepsilon - \varepsilon^2}$ and $p(-1+\varepsilon) = \sqrt{2\varepsilon - \varepsilon^2}$. It follows that if $u(x)$ and $v(x)$ satisfy the boundary conditions (iii), then

$$\lim_{\varepsilon \rightarrow 0^+} p(1-\varepsilon)[u'(1-\varepsilon)v(1-\varepsilon) - u(1-\varepsilon)v'(1-\varepsilon)] = 0$$

and

$$\lim_{\varepsilon \rightarrow 0^+} p(-1+\varepsilon)[u'(-1+\varepsilon)v(-1+\varepsilon) - u(-1+\varepsilon)v'(-1+\varepsilon)] = 0.$$

Therefore Eq. (17) is satisfied and the boundary value problem is *self-adjoint*.

(c). For $n \neq 0$,

$$\begin{aligned} n^2 \int_{-1}^1 \frac{T_0(x) T_n(x)}{\sqrt{1-x^2}} dx &= \int_{-1}^1 T_0(x) L[T_n] dx \\ &= \int_{-1}^1 L[T_0] T_n(x) dx \\ &= 0, \end{aligned}$$

since $L[T_0] = 0 \cdot T_0 = 0$. Otherwise,

$$\begin{aligned}
 n^2 \int_{-1}^1 \frac{T_n(x) T_m(x)}{\sqrt{1-x^2}} dx &= \int_{-1}^1 L[T_n] T_m(x) dx \\
 &= \int_{-1}^1 T_n(x) L[T_m] dx \\
 &= m^2 \int_{-1}^1 \frac{T_n(x) T_m(x)}{\sqrt{1-x^2}} dx .
 \end{aligned}$$

Therefore

$$(n^2 - m^2) \int_{-1}^1 \frac{T_n(x) T_m(x)}{\sqrt{1-x^2}} dx = 0 .$$

So for $n \neq m$,

$$\int_{-1}^1 \frac{T_n(x) T_m(x)}{\sqrt{1-x^2}} dx = 0 .$$