

**Problems and Solutions Section 6.7 (6.53 through 6.63)**

- 6.53** Calculate the response of Example 6.7.1 for  $l = 1$  m,  $E = 2.6 \times 10^{10}$  N/m<sup>2</sup> and  $\rho = 8.5 \times 10^3$  kg/m<sup>3</sup>. Plot the response using the first three modes at  $x = l/2$ ,  $l/4$ , and  $3l/4$ . How many modes are needed to represent accurately the response at the point  $x = l/2$ ?

**Solution:**

$$w(x, t) = \sum_{n=1}^{\infty} \left( \frac{0.02}{l^2 \sigma_n^2} (-1)^{n+1} \right) e^{-0.01 \omega_n t} \cos \omega_{2n} t \sin \sigma_n x$$

Where

$$\sigma_n = \frac{(2n-1)\pi}{2l}$$

$$\omega_n = \sigma_n \sqrt{\frac{E}{\rho}}$$

$$\omega_{dn} = 0.9999 \omega_n$$

For  $l = 1$  m

$$E = 2.6 \times 10^{10} \text{ N/m}^2$$

$$\rho = 8.5 \times 10^3 \text{ kg/m}^3$$

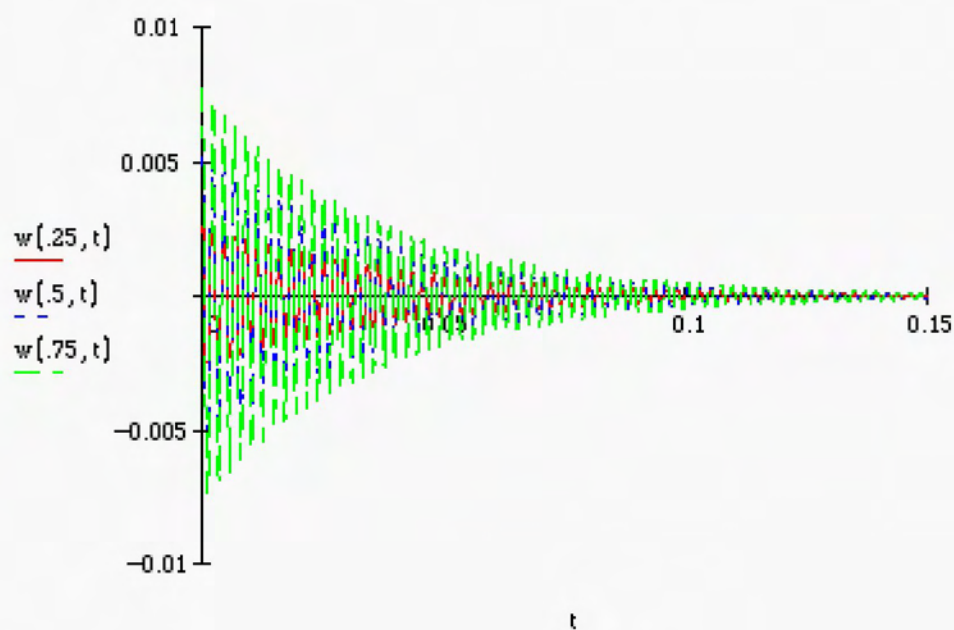
Response using first three modes at  $x = \frac{l}{2}, \frac{l}{4}, \frac{3l}{4}$  plotted below.

Three modes accurately represents the response at  $x = \frac{l}{2}$ . The error between a three and higher mode approximation is less than 0.2%.

$$\sigma(n) := \frac{2 \cdot n - 1}{2} \cdot \pi \quad E := 2.6 \cdot 10^{10} \quad \rho := 8.5 \cdot 10^3$$

$$\omega(n) := \sigma(n) \cdot \sqrt{\frac{E}{\rho}}$$

$$w(x, t) := \sum_{n=1}^3 \frac{.02}{\sigma(n)^2} (-1)^{n+1} \cdot e^{-0.01 \cdot \omega(n) \cdot t} \cdot \cos(\omega(n) \cdot t) \cdot \sin(\sigma(n) \cdot x)$$



**6.54** Repeat Example 6.7.1 for a modal damping ratio of  $\zeta_n = 0.01$ .

**Solution:** Using  $\zeta_n = 0.01$  and the frequency given in the example

$$\omega_{dn} = \omega_n \sqrt{1 - \zeta_n^2} = 0.995\omega_n, \quad \omega_n = \frac{2n-1}{2l} \sqrt{\frac{E}{\rho}} = \sigma_n \sqrt{\frac{E}{\rho}}$$

The time response is then  $T_n(t) = A_n e^{-0.1\omega_n t} \sin(\omega_{dn} t + \phi_n)$  and the total solution is:

$$w(x, t) = \sum_{n=1}^{\infty} A_n e^{-0.1\omega_n t} \sin(\omega_{dn} t + \phi_n) \sin \frac{(2n-1)}{2l} \pi x$$

The initial conditions are:

$$w(x, 0) = 0.01 \frac{x}{l} \text{ m and } w_t(x, 0) = 0$$

Therefore:

$$0.01 \frac{x}{l} = A_n \sin \phi_n \sin \sigma_n x$$

Multiply by  $\sin \sigma_m x$  and integrate over the length of the bar to get

$$0.01 \frac{(-1)^{m+1}}{l \sigma_m^2} = A_m \sin \phi_m \frac{1}{2} \quad m = 1, 2, 3, \dots$$

From the velocity initial condition

$$w_t(x, 0) = 0 = \sum_{n=1}^{\infty} A_n \left[ -0.1\omega_n \sin \phi_n + \omega_{dn} \cos \phi_n \right] \sin \sigma_n x$$

Again, multiply by  $\sin \sigma_m x$  and integrate over the length of the bar to get

$$A_m (-0.1\omega_n \sin \phi_n + \omega_{dn} \cos \phi_n) \frac{1}{2} = 0$$

Since  $A_m$  is not zero this yields:

$$\tan \phi_n = \frac{\sin \phi_n}{\cos \phi_n} = \frac{\sqrt{1 - \zeta_n^3}}{0.1} = 9.9499 \Rightarrow \phi_n = 1.4706 \text{ rad} = 84.3^\circ$$

Substitution into the equation from the displacement initial condition yields:

$$A_m = \frac{0.01}{l^2 \sigma_m^2} (-1)^{m+1} \frac{1}{\sin \phi_n} = \frac{0.0201}{l^2 \sigma_m^2} (-1)^{m+1}$$

The solution is then

$$w(x, t) = \sum_{n=1}^{\infty} \frac{0.01}{l^2 \sigma_m^2} (-1)^{m+1} e^{-0.1\omega_n t} \sin(\omega_{dn} t + \phi_n) \sin \frac{(2n-1)}{2l} \pi x$$

- 6.55** Repeat Problem 6.53 for the case of Problem 6.54. Does it take more or fewer modes to accurately represent the response at  $l/2$ ?

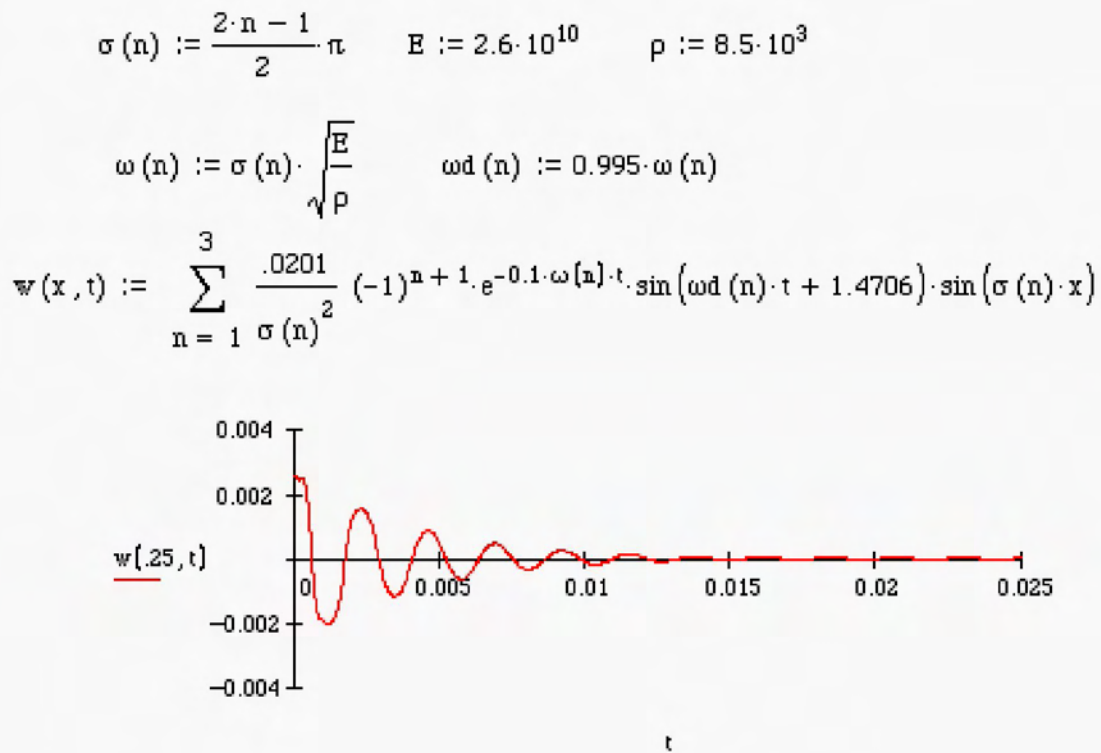
**Solution:** Use the result given in 6.54 and

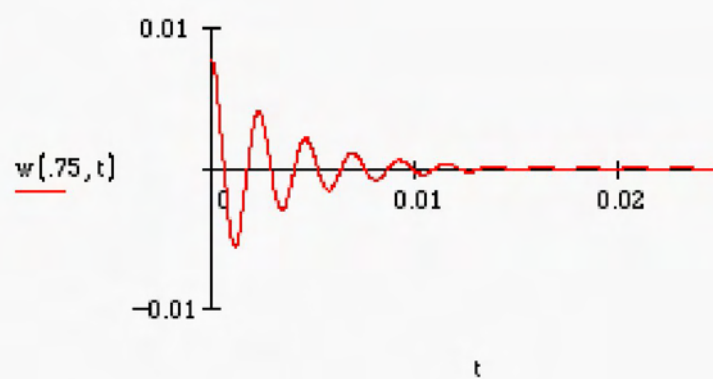
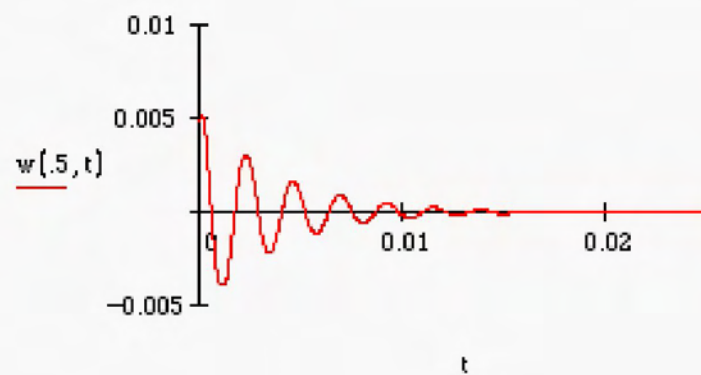
$$l = 1 \text{ m}$$

$$E = 2.6 \times 10^{10} \text{ N/m}^2$$

$$\rho = 8.5 \times 10^3 \text{ kg/m}^3$$

The response is plotted below at  $x = \frac{l}{4}, \frac{l}{2}, \frac{3l}{4}$ . An accurate representation of the response is obtained with three modes. The error between a three mode and a higher mode representation is always less than 0.2%. The results here are from Mathcad:





- 6.56** Calculate the form of modal damping ratios for the clamped string of equation (6.151) and the clamped membrane of equation (6.152).

**Solution:**

(a) For the string:

$$\begin{aligned}\rho w_{tt} + \gamma w_t - \tau w_{xx} &= 0 \\ \rho \phi_{tt} + \gamma \phi_t - \tau \phi'' &= 0 \\ \frac{\rho}{\tau} \frac{\phi_{tt}}{q} + \frac{\gamma}{\tau} \frac{\phi_t}{q} &= \frac{\phi''}{\phi} = -\sigma^2 \\ \phi_{tt} + \left(\frac{\gamma}{\rho}\right) \phi_t + \left(\frac{\tau}{\rho}\right) \sigma^2 \phi &= 0\end{aligned}$$

$\phi'' + \sigma^2 \phi = 0$  which has the solution  $\phi = A \sin \sigma x + B \cos \sigma x$ . The boundary conditions  $\phi(0) = \phi(l) = 0$  yield  $\sigma_n = \frac{n\pi}{l}$ ,  $n = 1, 2, 3, \dots$

$$\begin{aligned}\omega_n^2 &= \left(\frac{\tau}{\rho}\right) \sigma_n^2 = \frac{\tau}{\rho} \left(\frac{n\pi}{l}\right)^2 \\ 2\zeta_n \omega_n &= \frac{\gamma}{\rho} \\ \zeta_n &= \frac{\gamma}{2\rho} \sqrt{\frac{\rho}{\tau}} \left(\frac{n\pi}{l}\right) \\ \zeta_n &= \frac{\gamma}{2\sqrt{\rho\tau}} \left(\frac{n\pi}{l}\right)\end{aligned}$$

(b) For the membrane

$$\begin{aligned}\frac{\rho}{\tau} w_{tt} + \frac{\gamma}{\tau} w_t &= w_{xx} + w_{yy} \\ \left(\frac{\rho}{\tau}\right) XY \phi_{tt} + \left(\frac{\gamma}{\tau}\right) XY \phi_t &= X''Yq + XY''q \\ \left(\frac{\rho}{\tau}\right) \frac{\phi_{tt}}{q} + \left(\frac{\gamma}{\tau}\right) \frac{\phi_t}{q} &= \frac{X''}{X} + \frac{Y''}{Y} = -\beta^2 \\ \phi_{tt} + \left(\frac{\gamma}{\rho}\right) \phi_t + \left(\frac{\tau}{\rho}\right) \beta^2 \phi &= 0\end{aligned}$$

$\frac{X''}{X} = -\frac{Y''}{Y} - \beta^2 = -\alpha^2$ . The boundary conditions are  $X(0) = X(l) = 0$  and  $Y(0) = Y(l) = 0$ . The two spatial solutions become

$$\begin{aligned} X'' + \alpha^2 X &= 0 & Y'' + \gamma^2 Y &= 0 \\ X &= A \sin \alpha x + B \cos \alpha x & Y &= C \sin \gamma x + D \cos \gamma x \\ B &= 0 & D &= 0 \\ \alpha_n &= \frac{n\pi}{l} \quad n = 1, 2, 3, \dots & \gamma_m &= \frac{m\pi}{l} \quad m = 1, 2, 3, \dots \end{aligned}$$

Thus

$$\begin{aligned} \beta_{mn}^2 &= (n^2 + m^2) \left( \frac{\pi}{l} \right)^2 \\ \omega_{mn}^2 &= \frac{\tau}{\rho} (n^2 + m^2) \left( \frac{\pi}{l} \right)^2 \\ 2\zeta_m n \omega_m n &= \frac{\gamma}{\rho} \\ \zeta_{mn} &= \frac{\gamma}{2\rho \omega_{mn}} = \frac{\gamma}{2\rho} \frac{1}{\sqrt{\frac{\tau}{\rho} (n^2 + m^2)}} \frac{l}{\pi} \\ \zeta_{mn} &= \frac{\gamma l}{2\sqrt{\rho \tau} (n^2 + m^2)} \end{aligned}$$

**6.57** Calculate the units on  $\gamma$  and  $\beta$  in equation (6.153).

**Solution:** The units are found from

$$\begin{aligned} \frac{\text{mg}}{\text{m}^3} \text{m}^2 \frac{\text{m}}{\text{s}^2} &= \gamma \frac{\text{m}}{\text{s}} \\ \frac{\text{kg}}{\text{s}^2} \frac{\text{s}}{\text{m}} &= \gamma \\ \gamma &= \frac{\text{kg}}{\text{m} \bullet \text{s}} \end{aligned}$$

- 6.58** Assume that  $E$ ,  $I$ , and  $\rho$  are constant in equations (6.153) and (6.154) and calculate the form of the modal damping ratio  $\zeta_n$ .

**Solution:**

If  $E$ ,  $I$ , and  $\rho$  are constant in equation 6.153 and 6.154. Then separation of variables works and the mode shapes become those given in table 6.4, which can be normalized so that  $\int_0^1 X_n X_m dx = \delta_{nm}$ . Substitution of  $w(x, t) = a_n(t)X_n(x)$  into equation (6.153) multiplying by  $X_m(x)$  and integrating over  $x$  yield the  $m$ th modal equation:

$$\rho A \ddot{a}_n(t) + \gamma \dot{a}_n(t) + \beta I \left( \frac{\omega_n^2}{c^2} \right) a_n(t) + EI \frac{\omega_n^2}{c^2} a_n(t) = 0$$

where equation (6.93) has been used to evaluate  $X''''$  and  $c^2 = EI / \rho A$ . Dividing by  $\rho A$  yields

$$\ddot{a}_n(t) + \left( \frac{\gamma}{\rho A} + \frac{\beta}{E} \omega_n^2 \right) \dot{a}_n(t) + \omega_n^2 a_n(t) = 0$$

which is the *sdo* form of windows 6.4. Thus the coefficients of must be and hence

$$2\zeta_n \omega_n = \frac{\gamma}{\rho A} + \frac{\beta}{E} \omega_n^2$$

and

$$\omega_n = \beta_n \sqrt{\frac{EI}{\rho A}}$$

$$\zeta_n = \frac{\gamma}{2\rho A \omega_n} + \frac{\beta}{E} \omega_n$$

where  $\beta_n$  are given in table 6.4.



**6.59** Calculate the form of the solution  $w(x,t)$  for the system of Problem 6.58.

**Solution:**

The form of the solution of the m time equation is just

$$A_n e^{-\zeta_n \omega_n t} \sin(\omega_{dn} t + \phi_n)$$

where  $\zeta_n$  and  $\omega_n$  are as given in problem 6.58,  $\omega_{dn} = \omega_n \sqrt{1 - \zeta_n^2}$ , and  $A_n$  and  $\phi_n$  are constants determined by initial conditions. The total solution is of the form

$$w(x,t) = \sum_{n=1}^{\infty} A_n e^{-\zeta_n \omega_n t} \sin(\omega_{dn} t + \phi_n) X_n(x)$$

where  $X_n(t)$  are the eigenfunctions given in table 6.4.

- 6.60** For a given cantilevered composite beam, the following values have been measured for bending vibration:

$$\begin{aligned} E &= 2.71 \times 10^{10} \text{ N/m}^2 & \rho &= 1710 \text{ kg/m}^3 \\ A &= 0.597 \times 10^{-3} \text{ m}^2 & l &= 1 \text{ m} \\ I &= 1.64 \times 10^{-9} \text{ m}^4 & \gamma &= 1.75 \text{ N s/m}^2 \\ \beta &= 20,500 \text{ N s/m}^2 \end{aligned}$$

Calculate the solution for the beam to an initial displacement of  $w_t(x,0) = 0$  and  $w(x,0) = 3 \sin \pi x$ .

**Solution:**

Using the values given and the formulas for  $a_n(t)$  from problem 6.58 the temporal equation becomes

$$\ddot{a}_n + (1.714 \times 10^{-6} \omega_n^2) a_n + \omega_n^2 a_n = 0$$

from problem 6.59,

$$w_t(x,t) \Big|_{t=0} = 0 = \sum A_n \left[ (-\zeta_n \omega_n) \sin \phi_n + \omega_{dn} \cos \phi_n \right] X_n(x)$$

and

$$w(x,0) = 3 \sin \pi x = \sum A_n \sin \phi_n X_n(x)$$

Multiplying by  $X_n(x)$  and integrating yields that

$$\zeta_n \omega_n \sin \phi_n = \omega_{dn} \cos \phi_n \quad \text{or} \quad \tan \phi_n = \frac{\omega_{dn}}{\zeta_n \omega_n}$$

and  $3 \int_0^l \sin \pi x X_n(x) dx = A_n \sin \phi_n$  so that

$$A_n = \frac{3 \int_0^l \sin \pi x X_n(x) dx}{\sin \phi_n} = \frac{3}{\sqrt{1 - \zeta_n^2}} \int_0^l \sin \pi x X_n(x) dx$$

where  $X_n(x)$  is given in table 6.4.

- 6.61** Plot the solution of Example 6.7.2 for the case  $w_t(x,0) = 0$ ,  $w(x,0) = \sin(n\pi x/\ell)$ ,  $\gamma = 10$  Ns/m<sup>2</sup>,  $\tau = 10^4$  N,  $\ell = 1$  m and  $\rho = 0.01$  kg/m<sup>3</sup>.

**Solution:** From equation (6.156) and the values given,  $\zeta_1 = 0.159/n$  or  $\zeta_n \omega_n = 500$  and  $\omega_{dn} = \sqrt{1 - 0.159^2}$ , so that:

$$w(x,t) = \sum_{n=1}^{\infty} A_n e^{-500t} \sin(\omega_{dn} t + \phi_n) \sin n\pi x$$

Applying the initial conditions yields

$$\int_0^1 \sin n\pi x \sin m\pi x dx = \sum_{n=1}^{\infty} A_n \sin(\phi_n) \int_0^1 \sin m\pi x \sin n\pi x dx$$

So that  $A_n \sin \phi_n = 0$  for all  $n$  except  $n = 1$ , and  $A_1 \sin \phi_1 = 1$ . So either  $\phi_n = 0$  or  $A_n = 0$

for  $n$  not zero. The other initial condition yields that  $\phi_n = \tan^{-1}(\frac{-\sqrt{1-\zeta_n^2}}{\zeta_n})$  so that

$A_n = 0$  for  $n$  not zero. Thus the system is only excited in the first mode. Then

$$\begin{aligned} w(x,t) &= A_1 e^{-500t} \sin(\omega_1 \sqrt{1-\zeta_1^2} t + \phi_1) \sin \pi x \\ &= -1.001 e^{-500t} \sin(3137.7t - 1.50) \sin \pi x \end{aligned}$$

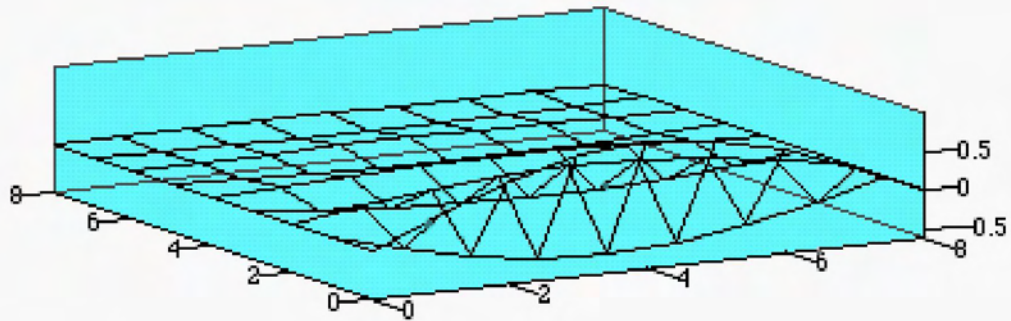
This is plotted in Mathcad below:

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N := 8    i := 0..N    j := 0..N    x_i := i * 1/8    t_j := j * 0.001
    ζ := 0.159    ω := π * 10^3    φ := -1.52    ωd := ω * sqrt(1 - ζ^2)
w(x,t) := -1.001 * e^(-ζ * ω * t) * (sin(π * x)) * sin(ωd * t + φ)

M_{i,j} := w(x_i, t_j)

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M

- 6.62** Calculate the orthogonality condition for the system of Example 6.7.2. Then calculate the form of the temporal solution.

**Solution:** Problem is to fill in the details of example 6.7.2 by checking the coefficients. Equation (6.155) by performing the integration.

- 6.63** Calculate the form of modal damping for the longitudinal vibration of the beam of Figure 6.14 with boundary conditions specified by equation (6.157).

**Solution:** This is a discussion problem. The boundary condition given in equation (6.157)

$$AEw_x(0,t) = kw(0,t) + c \frac{\partial w(0,t)}{\partial t}$$

$$AEw_x(l,t) = -kw(l,t) - c \frac{\partial w(l,t)}{\partial t}$$

Do not conform readily to separation of variables and lead to time dependent boundary conditions. However one approach is to treat the damper as applied forces of the bar  $cw_t(0,t)$  and  $-cw_t(l,t)$ . Following this approach the boundary conditions become

$$AEX'(0) = kX(0) \text{ and } AEX'(l) = -kX(l)$$

The general solution of the spatial equation of a bar has the form

$$X(x) = a \sin(\sigma x + b)$$

Where  $\sigma$  is the usual separation constant and  $a$  and  $b$  are constants. The first boundary condition yields that  $\phi = \tan^{-1}(AE/k)$ . The second boundary condition yields the characteristic equation

$$-(AE/k)\sigma_n = \tan(\sigma_n l + \phi)$$

Which can be solved for  $\sigma_n$  numerically. Note that  $\sigma_n$  are distinct so that from problem 6.39 the eigenfunctions are orthogonal, i.e. an can be calculated such that

$$X_n(x) = a_n \sin(\sigma_n x + \phi)$$

Are orthonormal. Following the procedure of example 6.8.11, the temporal solution for the forced response is

$$\begin{aligned} \ddot{w}_n(t) + \omega_n^2 w_n(t) &= \int_0^l [cw_t(0,t) - cw_t(l,t)] X_n(x) dx \\ &= \left\{ \int_0^l [cX_n(0) - cX_n(l)] X_n(x) dx \right\} \ddot{w}_n(t) \end{aligned}$$

Bring the  $\ddot{w}_n$  term to the left side and comparing its coefficient to  $2\zeta_n \omega_n$  yields

$$2\zeta_n \omega_n = \left\{ \int_0^l c [X_n(l) - X_n(0)] X_n(x) dx \right\}$$

The form of the modal damping ratio is thus

$$\zeta_n = \frac{ca_n^2}{2\omega_n \sigma_n} [\cos(\sigma_n l + \phi) - \cos \phi]$$

Where  $a_n^2$  is the normalization factor,  $\sigma_n$  are the eigenvalues  $\omega_n^2 = c^2 \sigma_n^2$  and  $\tan^{-1}(AE / k)$ .